

Propagation of chaos results for some McKean-Vlasov models

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Nonlinear processes and their applications, July 2019

Models of interest:

- Prototypical McKean-Vlasov dynamic:

$$dX_t = \left(\int b(X_t - x) \mu(t, dx) \right) dt + \sigma dW_t, \quad \mu(t) = \mathcal{L}(X_t);$$

- Langevin McKean-Vlasov dynamic:

$$\begin{cases} dY_t = V_t dt, \\ dV_t = \left(\int b(Y_t - y, V_t - v) \mu(t, dy, dv) \right) dt + \sigma dW_t. \end{cases}$$

- Conditional McKean Lagrangian model:

$$\begin{cases} dY_t = V_t dt, \\ dV_t = \mathbb{E}[b(V_t) | Y_t] dt + \sigma dW_t. \end{cases}$$

Generically, σ is a positive scalar or $\sigma \in \mathbb{R}^{d \times d}$ positive definite, and b measurable bounded.

Toy models (Sznitman 1991): Given $0 < T < \infty$ and $b : \mathbb{R}^d \rightarrow \mathbb{R}^d$ bounded Lipschitz continuous, $(X_0^i, (W_t^i; t \geq 0))$ independent copies of $(X_0, (W_t; t \geq 0))$,

$$X_t^{i,N} = X_0^i + \int_0^t \frac{1}{N} \sum_{j=1}^N b(X_s^{i,N} - X_s^{j,N}) ds + \sigma W_t^i,$$

$$X_t^{i,\infty} = X_0^i + \int_0^t \left(\int b(X_s^{i,\infty} - y) \mu(s, dy) \right) ds + \sigma W_t^i, \mu(t) = \mathcal{L}(X_t^{i,\infty}).$$

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$$X_t^{i,\infty} = X_0^i + \int_0^t \left(\int b(X_s^{i,\infty} - y) \mu(s, dy) \right) ds + \sigma W_t^i, \mu(t) = \mathcal{L}(X_t^{i,\infty}).$$

Theorem 1 (Sznitman 1989)

There exists $0 < C < \infty$, which doesn't depend on N , such that, for all $1 \leq i \leq N$,

$$W_1 \left(\mathcal{L}(X_t^{i,N}; t \leq T), \mathcal{L}(X_t^{i,\infty}; t \leq T) \right) \leq \mathbb{E} \left[\max_{0 \leq t \leq T} |X_t^{i,N} - X_t^{i,\infty}| \right] \leq \frac{C}{\sqrt{N}},$$

where W_1 is the (truncated) Monge-Kantorovich-Wasserstein distance:

$$W_1(P, Q) = \inf_{\pi \text{ coupling of } P \text{ and } Q} \int_{\mathcal{C}([0, T]; \mathbb{R}^d) \times \mathcal{C}([0, T]; \mathbb{R}^d)} (|x - y| \wedge 1) \pi(dx, dy).$$

Examples in the literature:

- Probabilistic interpretation of nonlinear pdes/stochastic particle methods :
 - Propagation of chaos for stochastic particle system with smooth or moderate interaction kernel: Funaki 1984; Oelschläger 1984, 1985; Léonard 1986; ... Méléard and Roelly-Coppoletta 1987; Méléard 1995; Bossy and Talay 1996; Jourdain and Méléard 1998; Antonelli and Kohastu-Higa 2002;
 - Burgers equation: Sznitman 1986; Bossy and Talay 1997; Jourdain 1997; ...
 - Incompressible Navier-Stokes equations: Marchioro and Pulvirenti 1982; Osada 1984; Méléard 2001; Fontbona 2004; ...
 - Conservative equations: Bossy and Jourdain 2000; Jourdain 2002; Jourdain and Reygner 2013; ...
 - Viscous Pressureless gas equation: Dermoune 2001; ...
 - Chemotaxis Keller-Segel model: Fournier and Jourdain 2017; Cattiaux and Pédèches 2017; Tomašević 2018; ...
- Mean Field Games and Optimal control problems for McKean-Vlasov dynamics (differential calculus on the space of probability measures): Lasry and Lions 2006; Lions "Cours au Collège de France" (2006-2012); Cardaliaguet *et al.* 2013; ...; Carmona and Delarue 2013, 2015, 2018; ...
- Recent quantitative "weak" propagation of chaos results: For $\psi : \mathbb{R}^d \rightarrow \mathbb{R}$ some test function,

$$\left| \mathbb{E}[\psi(X_T^{i,N})] - \mathbb{E}[\psi(X_T)] \right| \leq \frac{C}{N}.$$

Ref.: Kolokoltsov 2010; Kolokoltsov, Troeva and Wang 2014 ; Mischler and Mouhot 2013; Mischler, Mouhot and Wennberg 2014; Jourdain and Bencheikh 2019; Chassagneux, Szpruch and Tse 2019; Chaudru de Raynal and Frikha 2019, ...

Propagation of chaos result for general Vlasov models with non-smooth interaction kernel: Hauray and Jabin 2015, [Jabin and Wang 2016](#), 2017.

- **Interacting particle systems with bounded interacting kernel b :**

$$\begin{cases} Y_t^{i,N} = Y_0^{i,N} + \int_0^t V_s^{i,N} ds, \\ V_t^{i,N} = V_0^{i,N} + \frac{1}{N-1} \sum_{j=1, j \neq i}^N \int_0^t b(Y_s^{i,N} - Y_s^{j,N}) ds + \sigma W_t^i. \end{cases}$$

- **Mean field limit system:**

$$\begin{cases} Y_t = Y_0 + \int_0^t V_s ds, \\ Y_t = V_0 + \int_0^t \left(\int b(Y_s - x) \mu(s, dy) \right) ds + \sigma W_t, \\ \mathcal{L}(Y_t, V_t) = \mu(t). \end{cases}$$

Main approach: PDE analysis.

Propagation of chaos result for general Vlasov models with non-smooth interaction kernel: Hauray and Jabin 2015, [Jabin and Wang 2016](#), 2017.

• **Interacting particle systems with bounded interacting kernel b :** For $\mathbf{x}^N = (x_1, \dots, x_N) \in \mathbb{R}^d \times \dots \times \mathbb{R}^d$, $\mathbf{u}^N = (u_1, \dots, u_N) \in \mathbb{R}^d \times \dots \times \mathbb{R}^d$,

$$\begin{cases} \partial_t f^N + \mathbf{u}^N \cdot \nabla_{\mathbf{x}^N} f^N + \frac{1}{N-1} \sum_{j=1, j \neq i}^N b(x_i - x_j) \cdot \nabla_{u_i} f^N + \frac{\sigma^2}{2} \Delta_{\mathbf{u}^N} f^N = 0 \\ f^N(0, \mathbf{x}^N, \mathbf{u}^N) = f_0(\mathbf{x}^N, \mathbf{u}^N). \end{cases}$$

• **Related McKean-Vlasov model:** Related Fokker-Planck equation:

$$\begin{cases} \partial_t f + u \cdot \nabla_x f + \int b(x-y) f(t, y, v) dy dv \cdot \nabla_u f + \frac{\sigma^2}{2} \Delta_u f = 0, \\ f(0, x, u) = f_0(x, u). \end{cases}$$

• Propagation of chaos result in terms of the total variation distance:

$$\|\mu - \nu\|_{TV, \mathcal{P}(\mathbb{R}^d)} = \left| \sup_{A \in \mathcal{B}(\mathbb{R}^d)} (\mu(A) - \nu(A)) \right|, \mu, \nu \in \mathcal{P}(\mathbb{R}^d),$$

and the relative entropy (provided the density $d\mu/d\nu$ of μ w.r.t. ν exists):

$$H(\mu | \nu) = \int \log(d\mu/d\nu)(z) \mu(dz).$$

Propagation of chaos result: Define $f^{k,N}$ as the marginal p.d.f. of $(Y_t^{i,N}, V_t^{i,N})$, $i = 1, \dots, k$:

$$\begin{aligned} f^{k,N}(t, x_1, \dots, x_k, u_1, \dots, u_k) \\ = \int f^N(t, x_1, \dots, x_k, y_{k+1}, \dots, y_N, u_1, \dots, u_k, v_{k+1}, \dots, v_N) \end{aligned}$$

and $f^{k,\infty}$ as the marginal p.d.f. of (Y_t, V_t) , $i = 1, \dots, k$:

$$f^{k,\infty}(t, x_1, \dots, x_k, u_1, \dots, u_k) = f(t, x_1, u_1) \times \dots \times f(t, x_k, u_k).$$

Theorem 2 (Jabin and Wang 2016)

Assume independency of the initial positions, and that for some $\lambda > 0$,

$$\int_{(\mathbb{R}^d \times \mathbb{R}^d)} f_0(x, u) \left(|x|^{2\lambda} + |u|^{2\lambda} + \ln(f_0(x, u)) \right) < \infty.$$

Then

$$\forall t \geq 0, \|f^{k,N}(t) - f^{k,\infty}(t)\|_{L^1} \leq \sqrt{2H(f^{k,N}(t)|f^{k,\infty}(t))} \leq C/\sqrt{N}.$$

Element of proof:

- Preliminary steps: By Csiszár-Pinsker-Kullback inequality,

$$\|f^{k,N}(t) - f^{k,\infty}(t)\|_{L^1} \leq \sqrt{2H(f^{k,N}(t)|f^{k,\infty}(t))}.$$

Next, by the super-additive property of the renormalized relative entropy (see Hauray and Mischler 2014),

$$\frac{1}{k}H(f^{k,N}(t)|f^{k,\infty}(t)) \leq \frac{1}{N}H(f^{N,N}(t)|f^{N,\infty}(t)).$$

- Entropy estimate: Use the Fokker-Planck equation related to f^N to get

$$\begin{aligned} & \frac{d}{dt} \left(\frac{1}{N}H(f^{N,N}(t)|f^{N,\infty}(t)) \right) \\ & \leq \frac{1}{N}H(f^{N,N}(t)|f^{N,\infty}(t)) + \frac{1}{N} \int f^{N,\infty}(t, \mathbf{x}^N, \mathbf{u}^N) \exp\{|R_N(t, \mathbf{x}^N, \mathbf{u}^N)|\} d\mathbf{x}^N d\mathbf{u}^N \end{aligned}$$

where

$$R_N(t, \mathbf{y}^N, \mathbf{v}^N) = \frac{1}{N} \sum_{i,j=1}^N \left(\nabla_{v_i} \log(f^{N,\infty}) \cdot \left(b(y_i - y_j) - \int b(y_i - y) f(t, y, v) dv \right) \right)$$

The rest of the proof relies on the estimate

$\sup_N \int f^{N,\infty}(t, \mathbf{x}^N, \mathbf{u}^N) \exp\{|R_N(t, \mathbf{x}^N, \mathbf{u}^N)|\} d\mathbf{x}^N d\mathbf{u}^N < \infty$ and Gronwall's lemma.

- Jabin and Wang 2018: For $b \in W^{-1,\infty}$ (i.e. $b^{(k)}(x) = \sum_l \partial_{x_l} G^{k,l}(x)$, $G \in L^\infty$) on the d -dimensional torus \mathbb{T}^d :

$$X_t = \left(Z_0 + \int_0^t \left\{ F(Z_s) + \int b(Z_s - y) \mu(s, dy) \right\} ds + \sigma W_t \right) \text{ mod } 1, \quad X_t \sim \mu(t),$$

$$X_t^{i,N} = \left(X_0^i + \frac{1}{N} \sum_{j=1, j \neq i}^N \int_0^t \left\{ F(X_s^{i,N}) + b(X_s^{i,N} - X_s^{j,N}) \right\} ds + \sigma W_t \right) \text{ mod } 1,$$

Theorem 3 (Jabin and Wang 2018)

Assuming that $F, \nabla_x \cdot F$ are in L^∞ , $b, \nabla_x \cdot b \in W^{-1,\infty}$, that the Fokker-Planck equation related to $(X_t; t \leq T)$ admits a positive density in $L^\infty((0, T); W^{2,p}(\mathbb{T}^d))$ for all $p < \infty$ and that the initial entropy $H(f^N(0)|f^{N,\infty}(0))$ satisfies the following properties:

$$\sup_{N \rightarrow \infty} NH(f^N(0)|f^{N,\infty}(0)) < \infty.$$

Then

$$\text{ess sup}_{t \leq T} \|f^{k,N}(t) - f^{k,\infty}(t)\|_{L^1} \leq \frac{C}{\sqrt{N}}.$$

Starting point for a probabilistic analog: Bounded interaction and non-degenerated case: $\sigma \in \mathbb{R}^{d \times d}$ constant invertible:

$$X_t = X_0 + \int_0^t \int b(X_s - y) \mu(s, dy) ds + \sigma W_t, \quad X_t \sim \mu(t). \quad (1)$$

Corollary 4 (A corollary of Mishura and Veretennikov 2016)

Assume that $\mathbb{E}[|Z_0|^4] < \infty$, $b : \mathbb{R}^d \rightarrow \mathbb{R}^d$ is a bounded Borel function and $\sigma \in \mathbb{R}^{d \times d}$ is invertible. Then (1) admits at most one unique weak solution.

Elements of proof: Let

$$\widehat{X}_t = \widehat{X}_0 + \int_0^t \int b(\widehat{X}_s - y) \mu(s, dy) ds + \sigma \widehat{W}_t, \quad \widehat{X}_t \sim \mu(t).$$

be another solution to (1) and let P_T and \widehat{P}_T be the laws of $(X_t)_{0 \leq t \leq T}$ and $(\widehat{X}_t)_{0 \leq t \leq T}$. Then, again, by the Csiszar-Pinsker-Kullback inequality,

$$\|P_T - \widehat{P}_T\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))} \leq \sqrt{2H(P_T | \widehat{P}_T)} \left(= \sqrt{2\mathbb{E}_{P_T}[\log(dP_T/d\widehat{P}_T)]} \right).$$

As

$$dP_T/d\hat{P}_T = \exp \left\{ - \int_0^T \sigma^{-1} \left(b(X_t - y)\mu(t, dy) - \int b(X_t - y)\hat{\mu}(t, dy) \right) dW_t + \frac{1}{2} \int_0^T \left| \sigma^{-1} \left(b(X_t - y)\mu(t, dy) - \int b(X_t - y)\hat{\mu}(t, dy) \right) \right|^2 dt \right\},$$

where $(W_t; t \leq T)$ is a Brownian motion under P_T , it follows that

$$H(P_T | \hat{P}_T) = \frac{1}{2} \mathbb{E} \left[\int_0^T \left| \sigma^{-1} \left(\int b(X_t - y)\mu(t, dy) - \int b(X_t - y)\hat{\mu}(t, dy) \right) \right|^2 dt \right].$$

Since the above is bounded by $\|b\|_{L^\infty} \int_0^T \max_{s \leq t} \|\mu(s) - \hat{\mu}(s)\|^2 dt$, uniqueness follows by Gronwall's inequality.

Toward a propagation of chaos result: Instead of considering two solutions of a McKean-Vlasov dynamic, what would happen if we consider N independent copies of

$$dX_t^{i, \infty} = \left(\int b(X_t^{i, \infty} - y)\mu(t, dy) \right) dt + \sigma dW_t^i,$$

and the related interacting particle system:

$$dX_t^{i, N} = \frac{1}{N} \sum_{j=1}^N b(X_t^{i, N} - X_t^{j, N}) dt + \sigma dW_t^i ?$$

Answer for path-dependent McKean-Vlasov dynamics: Lacker 2018 (also Veretennikov April 2018 - Arxiv - for similar ideas on the queueing system GI/GI/1):

$$dX_t^{i,\infty} = B(t, (X_r^{i,\infty})_{0 \leq r \leq t}, \mathcal{L}((X_r^{i,\infty})_{0 \leq r \leq t})) dt + A(t, (X_r^{i,\infty})_{0 \leq r \leq t}) dW_t^i,$$

$$dX_t^{i,N} = B(t, (X_r^{i,N})_{0 \leq r \leq t}, \frac{1}{N} \sum_{i=1}^N \delta_{\{(X_r^{i,N})_{0 \leq r \leq t}\}}) dt + A(t, (X_r^{i,N})_{0 \leq r \leq t}) dW_t^i,$$

Theorem 5 (Propagation of chaos, Lacker 2018)

Assume that A is invertible and that the SDE $d\chi_t = A(t, (\chi_r)_{0 \leq r \leq t}) dW_t$, $\chi_0 \sim \mu_0$ has a unique strong solution. Assume also that $A^{-1}B$ is bounded,

$$|(A^{-1}B)(t, z, \mu) - (A^{-1}B)(t, z, \nu)| \leq K \|\mu - \nu\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))}$$

and: for all $\mu \in \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))$,

$$\nu \mapsto \int_{\mathcal{C}([0, T]; \mathbb{R}^d)} \int_0^T |A^{-1}(t, z) (B(t, z, \mu) - B(t, z, \nu))|^2 dt \nu(dz)$$

is sequentially continuous. Then the following propagation of chaos holds:

$$\lim_{N \rightarrow \infty} \|P^{k,N} - P^{k,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} = \lim_{N \rightarrow \infty} H(P^{k,\infty} | P^{k,N}) = 0,$$

for

$$P^{k,\infty} = \mathbb{P} \circ ((X_t^{1,\infty}, \dots, X_t^{k,\infty})_{t \leq T})^{-1}, \quad P^{k,N} = \mathbb{P} \circ ((X_t^{1,N}, \dots, X_t^{k,N})_{t \leq T})^{-1}.$$

Elements of proof for the TV and entropy limits: Preliminary, by a LDP, show that for all set $U \in \mathcal{B}(\mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d)))$ containing $\mathcal{L}((X_t)_{0 \leq t \leq T})$,

$$\lim_{N \rightarrow \infty} \mathbb{P}(\Gamma_t^N \notin U) = 0, \quad \Gamma_t^N := \frac{1}{N} \sum_{i=1}^N \delta_{\{(X_r^{i,N})_{0 \leq r \leq t}\}}.$$

Next, by Csiszár-Pinsker-Kullback inequality,

$$\|P^{k,N} - P^{k,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \leq \sqrt{2H(P^{k,\infty} | P^{k,N})}$$

where $H(P^{k,\infty} | P^{k,N})$ is the relative entropy linked to introduce the interacting kernel in the k first components of the McKean-Vlasov model. Using the exchangeability of the partially-interacting particle system, Girsanov's transformation and appropriate L^p -estimates on the density $dP^{k,\infty}/dP^{k,N}$,

$$H(P^{k,\infty} | P^{k,N}) \leq \frac{k \|A^{-1}B\|_{L^\infty} T e^{4kT \|A^{-1}B\|_{L^\infty}}}{2} \times \left(\mathbb{E} \left[\int_0^T \left| (A^{-1}B)(t, (X_s^{1,N})_{0 \leq s \leq t}, \Gamma_t^N) - (A^{-1}B)(t, (X_r^{1,N})_{0 \leq r \leq t}, \mathcal{L}((X_r^{1,\infty})_{0 \leq r \leq t})) \right|^2 dt \right] \right)^{1/2}$$

The regularity assumptions on A^{-1} and B ensures that the upper-bound tends to 0 as $N \rightarrow \infty$.

Toward a quantitative propagation of chaos result: Using the estimate

$$\|P^{k,N} - P^{k,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \leq \sqrt{2H(P^{k,N} | P^{k,\infty})}$$

and the super-additivity property of the relative entropy on $\mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))$:

$$H(P^{k,N} | P^{k,\infty}) = \frac{k}{N} H(P^{N,N} | P^{N,\infty}).$$

\Rightarrow Provided that $\max_{N \geq 1} H(P^{N,N} | P^{N,\infty}) < \infty$, an (optimal) quantitative propagation of chaos emerges.

Sufficient condition: For $\delta > 0$,

$$\mathbb{E} \left[(Z_T^{N,\infty})^{1+\delta} \right] < \infty,$$

where $Z_T^{N,\infty}$ is the exponential martingale allowing to pass from the N -McKean-Vlasov model to the N -interacting particle system:

$$Z_T^{N,\infty} := dP^{N,N} / dP^{N,\infty} = \exp \left\{ - \sum_{i=1}^N \int_0^T \Delta B_t^{i,N} \cdot dW_t^i - \frac{1}{2} \int_0^T \sum_{i=1}^N |\Delta B_t^{i,N}|^2 dt \right\},$$

$$\Delta B_t^{i,N} = A^{-1}(t, (X_r^{i,\infty})_{0 \leq r \leq t}) \times \left(B(t, (X_r^{i,\infty})_{0 \leq r \leq t}, \frac{1}{N} \sum_{j=1}^N \delta_{\{(X_r^{j,\infty})_{0 \leq r \leq t}\}}) - B(t, (X_r^{i,\infty})_{0 \leq r \leq t}, \mathcal{L}((X_r^{i,\infty})_{0 \leq r \leq t})) \right).$$

Alternative: Coming back to the definition of the total variation distance between $P^{k,\infty}$ and $P^{k,N}$:

$$\begin{aligned}
 & \|P^{k,N} - P^{k,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \\
 &= \sup_{A \in \mathcal{B}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \left| \mathbb{P}((X^{1,N}, \dots, X^{k,N}) \in A) - \mathbb{P}((X^{1,\infty}, \dots, X^{k,\infty}) \in A) \right| \\
 &= \sup_{A \in \mathcal{B}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \left| \mathbb{E}[Z_T^{N,\infty} \mathbb{1}_{\{(X^{1,\infty}, \dots, X^{k,\infty}) \in A\}}] - \mathbb{P}((X^{1,\infty}, \dots, X^{k,\infty}) \in A) \right| \\
 &= \sup_{A \in \mathcal{B}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \left| \mathbb{E} \left[\left(Z_T^{N,\infty} - 1 \right) \mathbb{1}_{\{(X^{1,\infty}, \dots, X^{k,\infty}) \in A\}} \right] \right| \\
 &= \sup_{A \in \mathcal{B}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \left| \mathbb{E} \left[\sum_{i=1}^N \int_0^T Z_t^{N,\infty} \Delta B_t^{i,N} dW_t^i \mathbb{1}_{\{(X^{1,\infty}, \dots, X^{k,\infty}) \in A\}} \right] \right| \\
 &\quad \text{(applying Itô's formula)} \\
 &= \sup_{A \in \mathcal{B}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \left| \mathbb{E} \left[\sum_{i=1}^k \int_0^T Z_t^{N,\infty} \Delta B_t^{i,N} dW_t^i \mathbb{1}_{\{(X^{1,\infty}, \dots, X^{k,\infty}) \in A\}} \right] \right| \\
 &\quad \text{(since } (X_t^{1,\infty}, W_t^1; t \leq T), \dots, (X_t^{N,\infty}, W_t^N; t \leq T) \text{ are i.i.d.)} \\
 &= \mathbb{E} \left[\left| \mathbb{E} \left[\sum_{i=1}^k \int_0^T Z_t^{N,\infty} \Delta B_t^{i,N} dW_t^i \mid (X^{1,\infty}, \dots, X^{k,\infty}) \right] \right| \right].
 \end{aligned}$$

An illustrative example: Taking the prototypical case: Diffusion A constant, b bounded,

$$X_t^{i,N} = X_0^i + \int_0^t \frac{1}{N} \sum_{j=1}^N b(X_s^{i,N} - X_s^{j,N}) ds + \sigma W_t^i,$$

$$X_t^{i,\infty} = X_0^i + \int_0^t \left(\int b(X_s^{i,\infty} - y) \mu(s, dy) \right) ds + \sigma W_t^i, \mu(t) = \mathcal{L}(X_t^{i,\infty}).$$

Then,

$$\Delta B_t^{i,N} = \sigma^{-1} \left(\frac{1}{N} \sum_{j=1}^N b(X_t^{i,\infty} - X_t^{j,\infty}) - \int b(X_t^{i,\infty} - y) \mu(t, dy) \right)$$

Rough estimate: For $k = 1$, using BDG and Hölder inequalities,

$$\begin{aligned} \|P^{k,N} - P^{k,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))} &= \mathbb{E} \left[\left| \mathbb{E} \left[\int_0^T Z_t^{N,\infty} \Delta B_t^{i,N} dW_t^i \mid (X^{1,\infty}, \dots, X^{k,\infty}) \right] \right| \right] \\ &\leq \sum_{i=1}^k \left(\mathbb{E} \left[\int_0^T (Z_t^{N,\infty})^4 dt \right] \right)^{1/4} \left(\mathbb{E} \left[\int_0^T |\Delta B_t^{i,N}|^4 dt \right] \right)^{1/4}. \end{aligned}$$

Since $(X_t^{1,\infty}; t \leq T), \dots, (X_t^{N,\infty}; t \leq T)$ are i.i.d. $\mathbb{E} [|\Delta B_t^{i,N}|^4] \leq \frac{C_T}{N^2}$ and, provided that $\mathbb{E}[(Z_T^{N,\infty})^4] < \infty$ one can deduce an explicit rate of convergence:

$$\|P^{k,N} - P^{k,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))} \leq C_T \sqrt{k/N}.$$

\Rightarrow A quantitative propagation of chaos in terms of the TV-distance requires moment control on $Z_T^{N,\infty}$.

Local control for $Z_T^{N,\infty}$ for the prototypical McKean-Vlasov model:

$$Z_T^{N,\infty} = \exp \left\{ - \sum_{i=1}^N \int_0^T \Delta B_t^{i,N} \cdot dW_t^i - \frac{1}{2} \int_0^T \sum_{i=1}^N |\Delta B_t^{i,N}|^2 dt \right\},$$

$$\Delta B_t^{i,N} = \sigma^{-1} \left(\frac{1}{N} \sum_{j=1}^N b(X_t^{i,\infty} - X_t^{j,\infty}) - \int b(X_t^{i,\infty} - y) \mu(t, dy) \right).$$

Proposition 6 (Local control for $Z^{N,\infty}$, J' 2019, [?])

For all $0 < T_0 < T < \infty$, $0 < \kappa < \infty$,

$$\sup_N \mathbb{E} \left[\exp \left\{ \kappa \sum_{i=1}^N \int_{T_0}^{T_0+\delta} \Delta B_t^{i,N} \cdot dW_t^i \right\} \right] < \infty,$$

whenever $\delta < (16^2 \kappa T d^2 \|\sigma^{-1} b\|_{L^\infty}^2)^{-1}$.

Elements of proof (inspired by Krylov and Rockner 2005): Exponential's expansion:

$$\mathbb{E} \left[\exp \left\{ \kappa \sum_{i=1}^N \int_{T_0}^{T_0+\delta} \Delta B_t^{i,N} \cdot dW_t^i \right\} \right] = \sum_l \frac{\kappa^l}{l!} \mathbb{E} \left[\left(\sum_{i=1}^N \int_{T_0}^{T_0+\delta} \Delta B_t^{i,N} \cdot dW_t^i \right)^l \right]$$

$$\leq 2 \sum_l \frac{(\kappa)^{2l}}{(2l)!} \left(1 + \mathbb{E} \left[\left(\sum_{i=1}^N \int_{T_0}^{T_0+\delta} \Delta B_t^{i,N} \cdot dW_t^i \right)^{2l} \right] \right).$$

$$\begin{aligned}
& \mathbb{E} \left[\left(\sum_{i=1}^N \int_{T_0}^{T_0+\delta} \Delta B_t^{i,N} \cdot dW_t^i \right)^{2l} \right] \\
& \leq 2^{2l} l! \mathbb{E} \left[\left(\sum_{i=1}^N \int_{T_0}^{T_0+\delta} |\Delta B_t^{i,N}|^2 dt \right)^l \right] \\
& \leq 2^{2l} l! N^l \delta^{l-1} \mathbb{E} \left[\int_{T_0}^{T_0+\delta} \left| \frac{1}{N} \sum_{j=1}^N \sigma^{-1} \left(b(X_t^{i,\infty} - X_t^{j,\infty}) - \int b(X_t^{i,\infty} - y) \mu(t, dy) \right) \right|^{2l} dt \right].
\end{aligned}$$

Corollary 7 (Bougeron et al. [2])

Let X_1, X_2, \dots, X_n be a sequence of i.i.d. r.v.s' such that a.s. $|X_1| \leq \bar{m} < \infty$. Then

$$\mathbb{E} \left[\left(\sum_{i=1}^n (X_i - \mathbb{E}[X_i]) \right)^{2q} \right] \leq q! (2n\bar{m}^2)^q.$$

Owing to the $\delta < \bar{\delta}$ for $\bar{\delta} = (16^2 \kappa T d^2 \|\sigma^{-1} b\|_{L^\infty}^2)^{-1}$. We recover a series of the form

$$\sum_{l=1}^{\infty} a_l, \quad a_l = \frac{(l)! l! C^l \delta^l}{(2l)!},$$

for C depending only on $\kappa T, \|\sigma^{-1} b\|_{L^\infty}$. Choosing δ small enough ensures that the sum is finite.

Theorem 8 ([?])

Assume that σ is definite positive and $b : (0, T) \times \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}^d$ is a bounded Borel function. Then

$$\|P^{k,N} - P^{k,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \leq \frac{C_{k,d,T, \|\sigma^{-1}b\|_{L^\infty}}}{\sqrt{N}}.$$

Elements of proof: Decomposing the interval $[0, T]$ into $\cup_{m=0}^{M-1} [m\delta, (m+1)\delta \wedge T]$, $M = \lfloor T/\delta \rfloor$: and, for each $0 \leq m \leq M$, define the family of processes $(Y_t^{1,N,m,\infty})_{t \leq T}$, $(Y_t^{N,N,m,\infty})_{t \leq T}$ as

$$Y_t^{i,N,m} = \begin{cases} X_0^i + \int_0^t \int b(Y_s^{i,N,m} - y) \mu^i(s, dy) ds + \sigma \widehat{W}_t^i, & 0 \leq t \leq m\delta, \mu^i(t) = \mathcal{L}(Y_t^i), \\ Y_{m\delta}^{i,N,m} + \frac{1}{N} \sum_{j=1}^N \int_{m\delta}^t b(Y_s^{i,N,m} - Y_s^{j,N,m}) ds + \sigma(W_t^i - W_{\delta m}^i), & m\delta < t \leq T \wedge \delta M. \end{cases}$$

This system corresponds to a family of N -independent McKean-Vlasov, that progressively interact with others on a time interval of the form $(m\delta, T]$.

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$$Y_t^{i,N,m} = \begin{cases} X_0^i + \int_0^t \int b(X_s^{i,\infty} - y) \mu(s, dy) ds + \sigma \widehat{W}_t^i, & 0 \leq t \leq m\delta, \\ X_{m\delta}^{i,\infty} + \frac{1}{N} \sum_{j=1}^N \int_{m\delta}^t b(X_s^{i,\infty} - X_s^{j,\infty}) ds + \sigma(W_t^i - W_{\delta m}^i), & m\delta < t \leq T \wedge \delta M. \end{cases}$$

This system corresponds to a family of N -independent McKean-Vlasov, that progressively interact with others on a time interval of the form $(m\delta, T]$.

By the triangular inequality

$$\|P^{1,\infty} - P^{1,N}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))} \leq \sum_{m=0}^M \|P^{1,m+1,N} - P^{1,m,N}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))}$$

where $P^{1,m+1,N}$ is the probability measure generated by $(Y_t^{i,N,m})_{0 \leq t \leq T}$. The exponential martingale which enables to change $Y^{i,m+1,N}$ into $Y^{i,m,N}$ is given by

$$\begin{aligned} & \exp \left\{ \sum_{i=1}^N \int_{m\delta}^{(m+1)\delta \wedge T} \left(\int b(Y_s^{i,N,m+1} - y) \mu^i(s, dy) - \frac{1}{N} \sum_{j=1}^N b(Y_s^{i,N,m} - Y_s^{j,N,m}) \right) dW_s^i \right. \\ & \quad \left. - \frac{1}{2} \sum_{i=1}^N \int_{m\delta}^{(m+1)\delta \wedge T} \left| \int b(Y_s^{i,N,m+1} - y) \mu^i(s, dy) - \frac{1}{N} \sum_{j=1}^N b(Y_s^{i,N,m} - Y_s^{j,N,m}) \right|^2 ds \right\} \\ & = \exp \left\{ - \sum_{i=1}^N \int_{m\delta}^{(m+1)\delta \wedge T} \Delta B_t^{i,N} \cdot dW_t^i - \frac{1}{2} \int_{m\delta}^{(m+1)\delta \wedge T} \sum_{i=1}^N |\Delta B_t^{i,N}|^2 dt \right\}. \end{aligned}$$

Each component $\|P^{1,m+1,N} - P^{1,m,N}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))}$ can then be controlled in terms of C/\sqrt{N} .

Remark: The proof only relies on the control of the exponential martingale which enables to pass from N -independent copies of the McKean-Vlasov models to the N -interacting particle systems and since this control is based on moment control of the difference $\Delta B_t^{i,N}$ given by some **concentration inequality/sub-gaussian moment controls** \Rightarrow Room for improvement.

- Langevin McKean-Vlasov models:

Theorem 9 ([?])

Assume that $b : \mathbb{R}^d \rightarrow \mathbb{R}^d$ is a bounded Borel function, σ definite positive. Then

$$\|\mathcal{L}(((Y_t^{i,N}, V_t^{i,N}); t \leq T)) - \mathcal{L}(((Y_t^{i,\infty}, V_t^{i,\infty}); t \leq T))\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^{kd}))} \leq C/\sqrt{N},$$

where

$$\begin{cases} dY_t^{i,N} = V_t^{i,N} dt, & 1 \leq i \leq N, \\ dV_t^{i,N} = \left(\frac{1}{N} \sum_{j=1}^N b(Y_t^{i,N} - Y_t^{j,N}, V_t^{i,N} - V_t^{j,N}) \right) dt + \sigma dW_t^i, \end{cases}$$

$$\begin{cases} dY_t^{i,\infty} = V_t^{i,\infty} dt, \\ dV_t^{i,\infty} = \left(\int b(Y_t^{i,\infty} - y, V_t^{i,\infty} - v) \mu(t, dy, dv) \right) dt + \sigma dW_t^i. \end{cases}$$

Rate of convergence for path-dependent McKean-Vlasov model: Consider the systems:

$$dX_t^{i,\infty} = B(t, (X_r^{i,\infty})_{0 \leq r \leq t}, \mathcal{L}((X_r^{i,\infty})_{0 \leq r \leq t})) dt + A(t, (X_r^{i,\infty})_{0 \leq r \leq t}) dW_t^i,$$

$$dX_t^{i,N} = B(t, (X_r^{i,N})_{0 \leq r \leq t}, \frac{1}{N} \sum_{i=1}^N \delta_{\{(X_r^{i,N})_{0 \leq r \leq t}\}}) dt + A(t, (X_r^{i,N})_{0 \leq r \leq t}) dW_t^i,$$

under the same assumptions of Lacker 2018. Assume further the **centering**

$$\mathbb{E}[B(t, (X_r^{i,\infty})_{0 \leq r \leq t}, \frac{1}{N} \sum_{i=1}^N \delta_{\{(X_r^{i,\infty})_{0 \leq r \leq t}\}})] = B(t, (X_r^{i,\infty})_{0 \leq r \leq t}, \mathcal{L}((X_r^{i,\infty})_{0 \leq r \leq t}))$$

Such condition ensures that

$$\mathbb{E} \left[\left(\int_0^T \left| (A^{-1}B)(t, (X_s^{1,\infty})_{0 \leq s \leq t}, \Gamma_t^N) - (A^{-1}B)(t, (X_r^{1,N})_{0 \leq r \leq t}, \mathcal{L}((X_r^{1,N})_{0 \leq r \leq t})) \right|^2 dt \right)^p \right] \leq \frac{p! \beta^p}{N^p}$$

and next

$$\|\mathcal{L}((X_t^{1,N}, \dots, X_t^{k,N}; t \leq T) - \mathcal{L}((X_t^{1,\infty}, \dots, X_t^{k,\infty}; t \leq T))\| - TV = \mathcal{O}(\sqrt{1/N}).$$

Note: More general cases would require more structural assumptions on B .

Conditional McKean Lagrangian model:

$$(*) \begin{cases} dX_t = U_t dt, \\ dU_t = \mathbb{E}[b(U_t) | X_t] dt + \sigma dW_t, \\ (X_0, U_0) \sim \mu_0 \in \mathcal{P}_2(\mathbb{R}^d \times \mathbb{R}^d). \end{cases}$$

Theorem 10 (Bossy, J. and Talay 2011)

Assume that $b : \mathbb{R}^d \rightarrow \mathbb{R}^d$ is bounded continuous, $\sigma \neq 0$ and μ_0 admits a Lebesgue density. Then there exists a unique solution to (*) and, as $N \rightarrow +\infty$ and $\epsilon \rightarrow 0^+$, the particle system

$$\begin{cases} X_t^{i,\epsilon,N} = X_0^i + \int_0^t U_s^{i,\epsilon,N} ds, \\ U_t^{i,\epsilon,N} = U_0^i + \int_0^t \frac{\frac{1}{N} \sum_{j=1}^N b(U_s^{j,\epsilon,N}) \phi_\epsilon(X_s^{i,\epsilon,N} - X_s^{j,\epsilon,N})}{\frac{1}{N} \sum_{j=1}^N \phi_\epsilon(X_s^{i,\epsilon,N} - X_s^{j,\epsilon,N}) + \epsilon} ds + \sigma W_t^i, \\ \{\phi_\epsilon\} \text{ family of smooth mollifiers,} \end{cases}$$

converges in the sense of propagation of chaos toward (*).

Rate of convergence: (joint work with S. Menozzi)

Intermediate smoothed dynamic:

$$X_t^\epsilon = X_0 + \int_0^t U_s^\epsilon ds, \quad U_t^\epsilon = U_0 + \int_0^t \frac{\mathbb{E}[b(\tilde{U}_s^\epsilon)\phi_\epsilon(x - \tilde{X}_s^\epsilon)]}{\mathbb{E}[\phi_\epsilon(x - \tilde{X}_s^\epsilon)] + \epsilon} \Big|_{x=X_s^\epsilon} ds + \sigma W_t,$$

for $(\tilde{X}_t, \tilde{U}_t)$ independent copy of (X_t, U_t) .







Error decomposition: For $P^{i,\epsilon,N}$, $P^{\epsilon,\infty}$ and P^∞ the respective law on $\mathcal{C}([0, T]; \mathbb{R}^d)$ of $(X^{i,\epsilon,N}, U^{i,\epsilon,N})$, (X^ϵ, U^ϵ) and (X, U) ,

$$\begin{aligned} & \|P^{i,\epsilon,N} - P^\infty\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))} \\ & \leq \underbrace{\|P^{i,\epsilon,N} - P^{\epsilon,\infty}\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))}}_{\sim C(\epsilon)/\sqrt{N} \text{ due to the singularity of the conditional expectation}} \\ & \quad + \underbrace{\|P^{\epsilon,\infty} - P^\infty\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))}}_{\text{requires a sharp estimate on the regularity of the marginal dist. of } P, P^\epsilon}. \end{aligned}$$

For the estimate on $\|P^{\epsilon,\infty} - P^\infty\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))}$: As

$$\begin{aligned} & \|P^{\epsilon,\infty} - P^\infty\|_{TV, \mathcal{P}(\mathcal{C}([0, T]; \mathbb{R}^d))} \leq \sqrt{2H(P^{\epsilon,\infty} | P^\infty)} \\ & \leq \sqrt{2\mathbb{E} \left[\int_0^T \left| \frac{\mathbb{E}[b(\tilde{U}_s^\epsilon)\phi_\epsilon(x - X_s^\epsilon)]}{\mathbb{E}[\phi_\epsilon(x - X_s^\epsilon)] + \epsilon} \Big|_{x=X_s} - \mathbb{E}[b(U_s) | X_s] \right|^2 ds \right]} \end{aligned}$$

the control of this expression relies on density estimates of $\mathcal{L}(X_t)$ and $\mathcal{L}(X_t^\epsilon)$.

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